## **Answers to exam-style questions**

## **Option D**

## $1 \checkmark = 1 \text{ mark}$

- 1 a i Luminosity is the total power radiated by a star.  $\checkmark$ 
  - ii Apparent brightness is the power received per unit area.  $\checkmark$

  - **b** They all fuse hydrogen into helium.  $\checkmark$  **c i** Using the mass-luminosity relation  $\frac{L_A}{L_2} = \left(\frac{M_A}{M_2}\right)^{3.5} \checkmark$

Hence 
$$L_{\rm A} = L_{\odot} \left(\frac{M_{\rm A}}{M_{\odot}}\right)^{3.5} = L_{\odot} \times 6.7^{3.5} = 778 \times 3.9 \times 10^{26} = 3.04 \times 10^{29} \approx 3.0 \times 10^{29} \, \text{W} \checkmark$$

ii From 
$$b = \frac{L}{4\pi d^2}$$
 we find  $d = \sqrt{\frac{L}{4\pi b}} = \sqrt{\frac{3.04 \times 10^{29}}{4\pi \times 1.7 \times 10^{-8}}} = 1.19 \times 10^{18} \text{ m} = \frac{1.19 \times 10^{18}}{3.09 \times 10^{16}} \text{ pc} = 38.5 \approx 38 \text{ pc}$ 

Hence 
$$p = \frac{1}{d} = \frac{1}{38.5} = 0.0260 \approx 0.026'' \checkmark$$
  
iii  $\frac{L_{\rm A}}{L_{\odot}} = 778 = \frac{\sigma 4\pi R_{\rm A}^2 T_{\rm A}^4}{\sigma 4\pi R_{\odot}^2 T_{\odot}^4} = \frac{R_{\rm A}^2 T_{\rm A}^4}{R_{\odot}^2 T_{\odot}^4} \checkmark$   
 $778 = \frac{R_{\rm A}^2}{R_{\odot}^2} \times 2.6^4 \checkmark$ 

Hence 
$$R_{\rm A} = R_{\odot} \frac{778}{2.6^4} \approx 17 R_{\odot}$$

The parallax method measures the position of a star two times six months apart.  $\checkmark$ d i The shift of the angular position of the star relative to the background of the distance stars.  $\checkmark$ Allows measurement of the parallax angle which is the angle subtended by the earth's orbit radius at the star. 🗸

The distance in pc is the reciprocal of the parallax angle in arc seconds.  $\checkmark$ 

ii Yes because it is larger than the limit of 0.01 arc seconds.  $\checkmark$ 

- 2 a Light reaching the Earth must go through the outer layers of the star.  $\checkmark$ Photons whose energy corresponds to differences in energy between energy levels of the atoms of the star may be absorbed and so will be missing in the received light.  $\checkmark$ Because the energy level differences are specific atoms determination of the chemical composition is then possible. 🗸
  - **b** i Type O stars are very hot and most of the hydrogen is ionised.  $\checkmark$ Hence hydrogen cannot absorb any photons. ✓
    - ii An M type star is relatively cool so that hydrogen atoms are mostly in their ground state.  $\checkmark$ And these can only absorb ultraviolet photons not visible light photons.  $\checkmark$
  - **c** The surface temperature/its magnetic field/its rotational speed.  $\checkmark$
- The surface of the star periodically expands and contracts.  $\checkmark$ 3 a It expands because radiation ionises helium atoms in the star's outer layers and the released electrons heat up the star expanding it. 🗸

When most of the helium is ionised, radiation leaves the star so the star cools and contracts.  $\checkmark$ 

There is a relation between the average luminosity of the star and the period of variation of the luminosity.  $\checkmark$ b So measuring the period gives the luminosity.  $\checkmark$ 

Measuring the (average) apparent brightness (by observing the star over a period of days) allows determination of the distance.  $\checkmark$ 

To measure the distance to a galaxy it must be determined whether a specific Cepheid star belongs to that galaxy. 🗸

c The average apparent brightness is estimated to be 2.8 × 10<sup>-8</sup> W m<sup>-2</sup>. ✓
A period of 55 days corresponds to an average luminosity of about 20000 solar luminosities. ✓

And so 
$$d = \sqrt{\frac{L}{4\pi b}} = \sqrt{\frac{20000 \times 3.9 \times 10^{26}}{4\pi \times 2.8 \times 10^{-8}}} \approx 5 \times 10^{18} \text{ m} \checkmark$$
  
**a i** Using the mass-luminosity relation  $\frac{L_A}{L_\odot} = \left(\frac{M_A}{M_\odot}\right)^{3.} \checkmark$   
Hence  $L_A = L_\odot \left(\frac{M_A}{M_\odot}\right)^{3.5} = L_\odot \times 20^{3.5} = 3.6 \times 10^4 \times 3.9 \times 10^{26} = 1.4 \times 10^{31} \text{ W} \checkmark$   
**ii**  $\frac{L_A}{L_\odot} = 3.6 \times 10^4 = \frac{\sigma 4\pi R^2 T^4}{\sigma 4\pi R_\odot^2 T_\odot^4} = 1.2^2 \times \frac{T^4}{T_\odot^4} \checkmark$   
Hence  $\frac{T}{T_\odot} = \sqrt[4]{\frac{3.6 \times 10^4}{1.2^2}} = 12.6 \approx 13 \checkmark$ 

**b i** The surface temperature decreases. ✓ And the radius increases. ✓

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- ii A type II supernova is the explosion of a massive star after it has entered the red supergiant phase while a type Ia supernova involves mass accretion onto a white dwarf. ✓
   So in this case we have a type II supernova. ✓
- iii The star will be a neutron star. ✓
   In which the neutron degeneracy pressure. ✓
   Balances the gravitational pressure in the star. ✓
- c Blue line as shown starting approximately at the correct point.  $\checkmark$



- 5 a i Strip from top left diagonally to bottom right. ✓
  ii Region in lower left. ✓
  - ii Region in lower left. V
  - iii Region to the right above main sequence. ✓iv Region joining main sequence to red giants. ✓

**b i** 
$$\frac{L_{\rm X}}{L_{\rm Y}} = \frac{100}{0.01} = 10^4 = \frac{\sigma 4\pi R_{\rm X}^2 T_{\rm X}^4}{\sigma 4\pi R_{\rm Y}^2 T_{\rm Y}^4} = \frac{R_{\rm X}^2}{R_{\rm Y}^2} \checkmark$$
  
Hence  $\frac{R_{\rm X}}{R_{\rm Y}} = 10^2 \checkmark$ 

ii 
$$\frac{L_{\rm X}}{L_{\rm Z}} = 1 = \frac{\sigma 4\pi R_{\rm X}^2 T_{\rm X}^4}{\sigma 4\pi R_{\rm Z}^2 T_{\rm Z}^4} = \frac{R_{\rm X}^2}{R_{\rm Z}^2} \times \left(\frac{12000}{3000}\right)^4 = \frac{R_{\rm X}^2}{R_{\rm Z}^2} \times 256 \checkmark$$
  
Hence  $\frac{R_{\rm X}}{R_{\rm Z}} = \sqrt{\frac{1}{256}} = \frac{1}{16} \checkmark$ 

**c** Using the mass-luminosity relation  $\frac{L_X}{L_{\odot}} = 100 = \left(\frac{M_X}{M_{\odot}}\right)$ 

Hence 
$$\frac{M_{\rm X}}{M_{\odot}} = 100^{1/3.5} = 3.7$$
  $\checkmark$ 

**d** i Line similar to blue line on HR diagram.  $\checkmark$ 



ii Electron degeneracy pressure.  $\checkmark$ 

Balances the gravitational pressure in the star.  $\checkmark$ 

iii It has to be less than the Chandrasekhar limit of 1.4 solar mass.  $\checkmark$ 

- $6\,$  a Distant galaxies appear to move away from us with a speed that is proportional to their distance.  $\checkmark$ 
  - b The received wavelength. ✓Is larger than the wavelength at emission. ✓
  - c On a large scale, the space between galaxies stretches. ✓

Thus the wavelength of light emitted from a distant galaxy also stretches so it is larger at reception.  $\checkmark$ 

**d** i 
$$z = \frac{\Delta \lambda}{\lambda_0} = \frac{780 - 656}{656} = 0.189 \checkmark$$
  
 $z = \frac{\nu}{c} \Longrightarrow \nu = 5.67 \times 10^7 \approx 5.7 \times 10^7 \text{ m s}^{-1} \checkmark$   
**ii**  $z = \frac{R}{R_0} - 1 = 0.189$ , i.e.  $\frac{R}{R_0} = 1.189 \checkmark$   
 $\frac{R_0}{R_0} = 0.84 \checkmark$ 

*R*  **iii** The data give a Hubble value constant of:  $v = Hd \Rightarrow H = \frac{5.67 \times 10^7}{920 \times 10^6 \times 3.09 \times 10^{16}} = 1.99 \times 10^{-18} \text{ s}^{-1} \checkmark$ Hence  $T = \frac{1}{H} = \frac{1}{1.99 \times 10^{-18}} = 5.0 \times 10^{17} \text{ s} \left( = \frac{5.0 \times 10^{17}}{365 \times 24 \times 3600} = 1.6 \times 10^{10} \text{ year} \right) \checkmark$  e To see whether the universe accelerates or decelerates in its expansion distant objects of large redshift had to be investigated. ✓

In order to establish the relation between distance and redshift.  $\checkmark$ 

Type Ia supernovae were chosen because their peak luminosity is known and hence the distance could be established by measuring the apparent brightness.  $\checkmark$ 

7 a According to the hot big bang model the early universe contained radiation at very high temperature. ✓
 As the universe expanded it cooled and the peak wavelength shifted to large microwave wavelengths with a black body spectrum. ✓

Which is what is being observed.  $\checkmark$ 

**b** i 
$$z = \frac{R}{R_0} - 1 = \frac{T_0}{T} - 1 \checkmark$$
  
 $z = \frac{3 \times 10^3}{2.7} - 1 = 1110 \checkmark$   
ii  $\frac{R}{R_0} - 1 = 1110 \checkmark$ 

Hence 
$$\frac{R_0}{R} = 9 \times 10^{-4}$$

- 8 a The result follows from  $\frac{GMm}{r^2} = m \frac{v^2}{r}$ .
  - **b** With M = kr the result in **a** becomes  $v = \sqrt{\frac{Gkr}{r}} = \sqrt{Gk}$  a constant.
  - c i The rotation curve becomes flat at large distances from the galactic centre. ✓
     This is consistent with a mass distribution as in b. ✓
     In other words with substantial mass far from the galactic centre. ✓
    - ii Small planets/brown dwarfs/black dwarfs. ✓ Neutrinos/exotic particle predicted by supersymmetry. ✓
- 9 a The Jeans criterion states that cloud of gas will begin to collapse under its own gravitation.  $\checkmark$ 
  - When the gravitational potential energy of the cloud exceeds the total random kinetic energy of its particles. ✓
    b Blue line as shown. ✓



**c i** Use M = Nm to eliminate N so that  $\frac{GM^2}{R} \approx \frac{3}{2} \frac{M}{m} kT$  or  $\frac{GM}{R} \approx \frac{3}{2} \frac{kT}{m} \checkmark$ Now eliminate the mass M through:  $M = \rho V = \rho \frac{4\pi R^3}{3}$  so that  $\frac{G\rho 4\pi R^3}{3R} \approx \frac{3}{2} \frac{kT}{m} \checkmark$ Cancelling powers of R and simplifying gives the result  $\left(R^2 \approx \frac{9}{8\pi} \frac{kT}{mG\rho}\right) \checkmark$  **ii**  $R^2 \approx \frac{9kT}{8\pi G\rho m} = \frac{9 \times 1.38 \times 10^{-23} \times 100}{8\pi \times 6.67 \times 10^{-11} \times 1.8 \times 10^{-19} \times 2.0 \times 10^{-27}} \checkmark$  $R \approx 1.4 \times 10^{17} \text{ m }\checkmark$ 

**10 a i**  $v = H_0 R \checkmark$ 

ii 
$$E = \frac{1}{2}mv^2 - \frac{GMm}{R} \checkmark$$
  
 $E = \frac{1}{2}mH_0^2R^2 - \frac{G\rho 4\pi R^3m}{3R} = \frac{1}{2}mH_0^2R^2 - \frac{G\rho 4\pi R^2m}{3}\checkmark$ 

Factoring gives the result.

iii To escape to infinity the total energy must be zero.  $\checkmark$ 

This means that 
$$H_0^2 - \frac{G\rho 4\pi}{3} = 0.$$

From which the result follows.

iv 
$$\rho = \frac{3 \times \left(\frac{68 \times 10^3}{10^6 \times 3.09 \times 10^{16}}\right)^2}{8\pi \times 6.67 \times 10^{-11}} \checkmark$$
  
 $\rho \approx 9 \times 10^{-27} \text{ kg m}^{-3} \checkmark$ 

**b** In cosmological models with matter density parameters  $\rho_{\rm m}$  and dark energy density  $\rho_{\Lambda}$  the significance of the critical density is that when  $\rho_{\rm m} + \rho_{\Lambda} = \rho_{\rm c}$ .

The geometry of the universe is flat, i.e. it has zero curvature.  $\checkmark$ 

11 a The CMB is very isotropic which means that we observe the same spectrum in every direction.  $\checkmark$  However there are small deviations from perfect isotropy in the sense that, in different directions, the temperature deviates from the mean temperature of T = 2.723 K by very small amounts (of order

$$\frac{\Delta T}{T} \approx 10^{-5}). \checkmark$$

- b These deviations are significant because fluctuations in temperature imply fluctuations in density. ✓
   And these are required if structures are to develop in the universe. ✓
   They are also significant because the magnitude of the fluctuations depends on the geometry of the universe. ✓
   Hence study of the fluctuations places limits on the geometry of the universe. ✓
- 12 a Elements are produced by nuclear fusion. ✓
   And nuclear fusion becomes energetically impossible past the peak of the binding energy per nucleon curve which is at iron. ✓
  - b These are produced mainly by neutron capture. ✓
    Nuclei may absorb neutrons which are abundant in a supernova. ✓
    As these decay by beta decay. ✓
    Nuclei with higher atomic number than iron are produced. ✓
  - c The CNO cycle requires the fusion of nuclei of carbon, nitrogen and oxygen and since these have a high atomic number the Coulomb barrier that must be overcome is larger than that for hydrogen and helium. ✓ And this requires higher temperatures that are found in the more massive stars. ✓